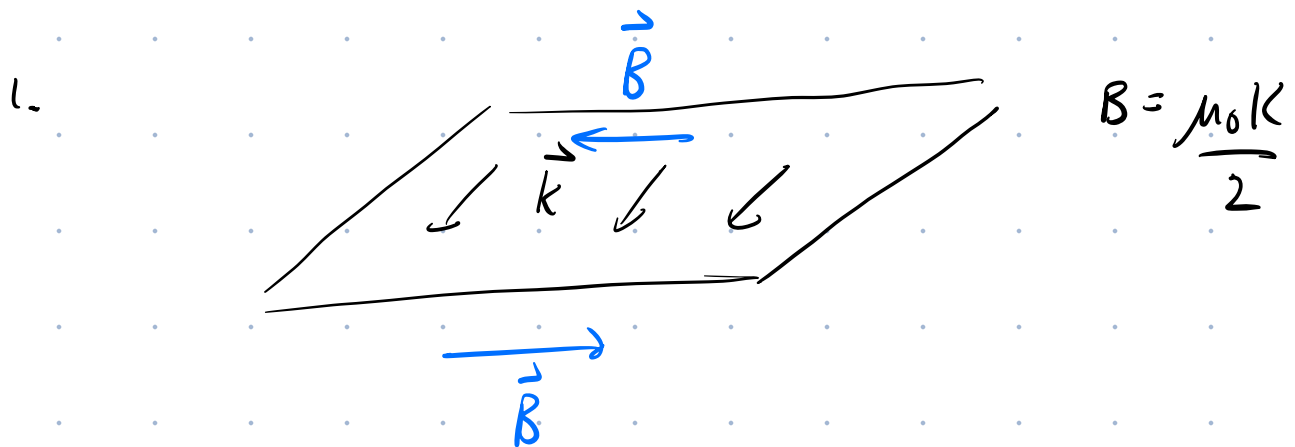
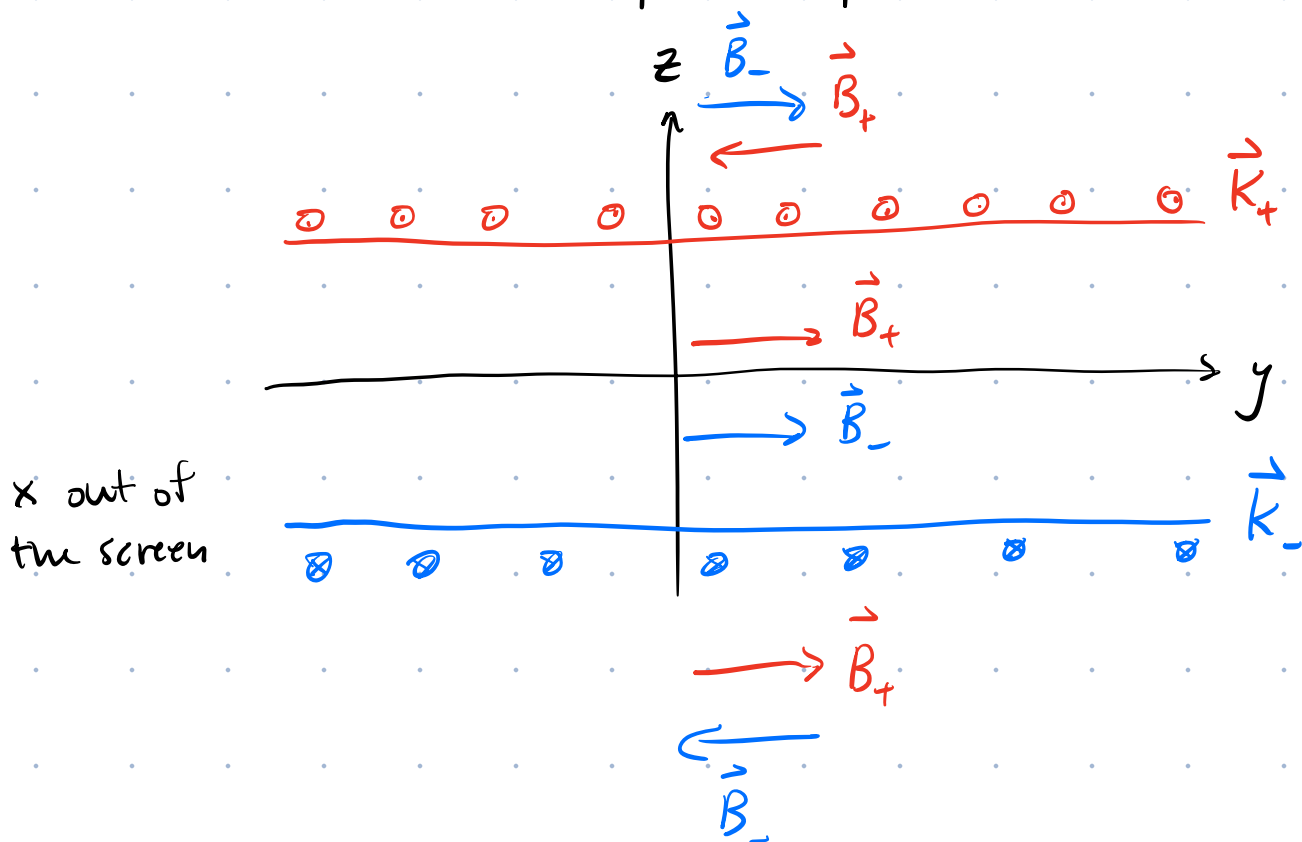




(a) Side View of a pair of parallel sheets.



Between the sheets, the $\vec{B}_+$ & $\vec{B}_-$ contributions add constructively s.t.

$$\boxed{\vec{B}_{in} = \mu_0 K \hat{y}} \quad (\text{for } -a < z < a)$$

Above and below, the two contribution cancel
s.t. $\vec{B}_{\text{out}} = 0$ (for $z > a$ and $z < -a$)

(b) Know that $\vec{B} = \vec{\nabla} \times \vec{A}$.

$\therefore$ Given $\vec{B}$, we can try to deduce $\vec{A}$,

For infinite sheets in x - y directions,
expect $\vec{A} = A(z)\hat{y}$.

Furthermore, we usually expect $\vec{A}$ to be parallel
or antiparallel to the current.

$$\therefore \vec{A} = A(z)\hat{y}$$

$$\therefore \vec{\nabla} \times \vec{A} = \begin{vmatrix} \hat{x} & \hat{y} & \hat{z} \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ A(z) & 0 & 0 \end{vmatrix} = \frac{dA}{dz} \hat{y}$$

$$\text{So we have } \frac{dA}{dz} = \begin{cases} 0 & z < -a \\ \mu_0 K & -a < z < a \\ 0 & z > a \end{cases}$$

(i) start w/ $-a < z < a$

$$dA = \mu_0 k dz \rightarrow \vec{A} = \mu_0 k z \hat{x}$$

Note that z changes sign when z passes through zero.

$$\text{Clearly } \vec{\nabla} \times \vec{A} = \mu_0 k \hat{y} = \vec{B} \text{ as required. } \checkmark$$

$$\text{It is also clear that } \vec{\nabla} \cdot \vec{A} = 0 \quad \checkmark$$

(ii) Next, consider $z > a$ s.t. $B = 0$.

In this case $\vec{\nabla} \times \vec{A} = 0$ and we require $\vec{A}$ to be continuous at $z = +a$

From (i) $\vec{A} = \mu_0 k a \hat{x}$ at $z = a$.

$$\therefore \text{ we take } \vec{A} = \mu_0 k a \hat{x} \quad \forall z \geq a.$$

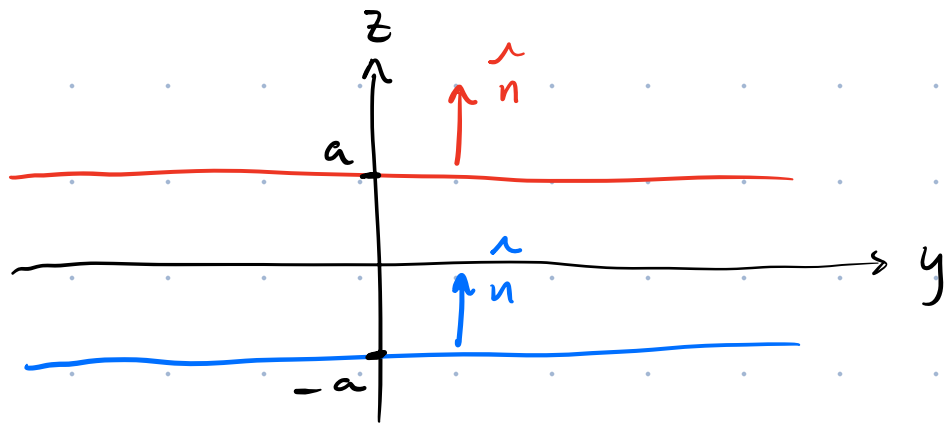
Likewise, by the same reasoning,

$$\vec{A} = -\mu_0 k a \hat{x} \quad \forall z \leq -a$$

(c) the other boundary condition on $\vec{A}$ is

$$\frac{\partial \vec{A}_{\text{above}}}{\partial n} - \frac{\partial \vec{A}_{\text{below}}}{\partial n} = -\mu_0 \vec{K}$$

where n is in the dir'n $\perp$ to $\vec{K}$. Take our normal vectors to point along $+\hat{z}$ dir'n.



Start w/ $z = +a$. $n = z$

$$\therefore \frac{\partial \vec{A}_{\text{above}}}{\partial z} - \frac{\partial \vec{A}_{\text{below}}}{\partial z} = -\mu_0 K \hat{x}$$

$$\Rightarrow \frac{\partial}{\partial z} (\cancel{\mu_0 K a \hat{x}}) - \underbrace{\frac{\partial}{\partial z} (\mu_0 K z \hat{x})}_{-\mu_0 K \hat{x}} = -\mu_0 K \hat{x} \quad \checkmark$$

Next, consider $z = -a$

$$n = \hat{z}$$

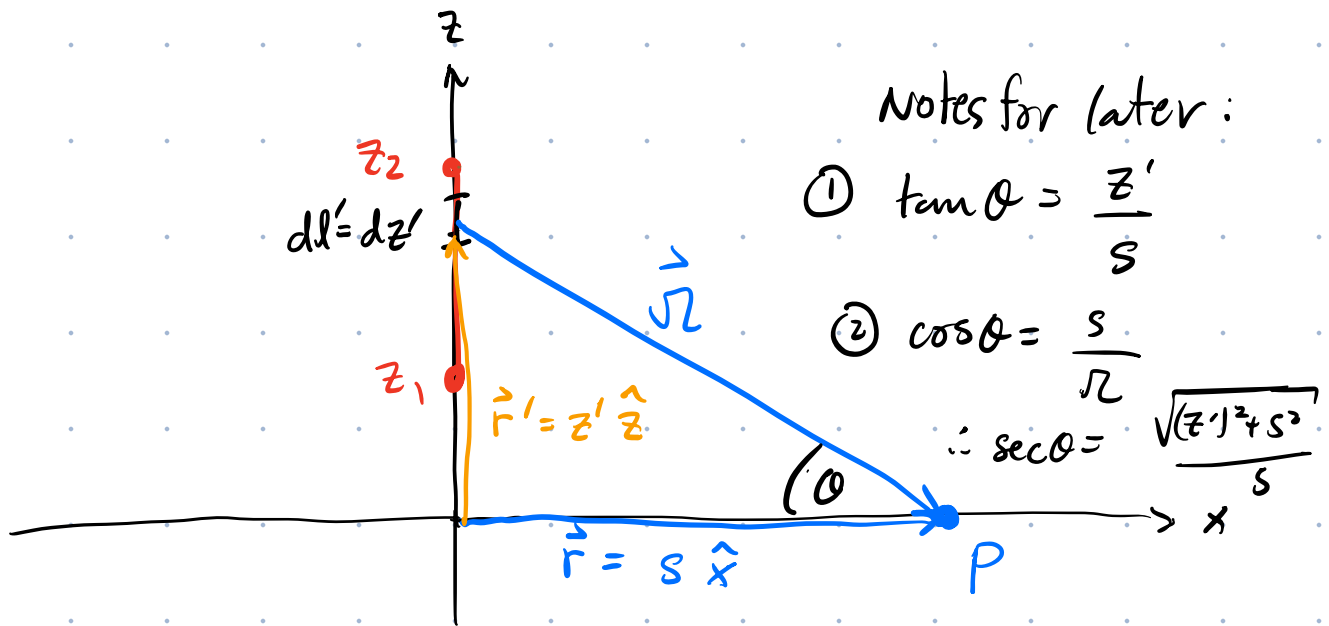
$$\frac{\partial \vec{A}_{\text{above}}}{\partial n} - \frac{\partial \vec{A}_{\text{below}}}{\partial n} = +\mu_0 K \hat{x}$$

b/c $\vec{K}$ is in the $-\hat{x}$ dir'n.

$$\underbrace{\frac{\partial}{\partial z} (\mu_0 K z \hat{x})}_{\mu_0 K \hat{x}} - \cancel{\frac{\partial}{\partial z} (-\mu_0 K a \hat{x})} = \mu_0 K \hat{x} \quad \checkmark$$

Note that "above" means the region of space that $\hat{n}$ points towards.

2.



We can arbitrarily place our "field pt" P in on the x-axis

Use cylindrical coords.

From the diagram $\vec{r} = s \hat{s}$
 $\vec{r}' = z' \hat{z}$

$$\therefore \vec{R} = \vec{r} - \vec{r}' = s \hat{s} - z' \hat{z}$$

$$\vec{I} = I \hat{z}$$

$$R = \sqrt{(z')^2 + s^2}$$

To calc. $\vec{A}$, we use $\vec{A} = \frac{\mu_0}{4\pi} \int \frac{\vec{I} dl'}{R}$

$$\vec{A} = \frac{\mu_0}{4\pi} \int_{z_1}^{z_2} \frac{I dz' \hat{z}}{\sqrt{(z')^2 + s^2}}$$

make the substitution $z' = s \tan \theta$

$$\therefore dz' = s \sec^2 \theta d\theta$$

$$\sqrt{(z')^2 + s^2} = s \sqrt{\tan^2 \theta + 1} = s \sec \theta$$

$$\therefore \vec{A} = \frac{\mu_0 I \hat{z}}{4\pi} \int_{\theta_1}^{\theta_2} \frac{\cancel{s} \sec^2 \theta d\theta}{\cancel{s} \cancel{\sec \theta}} = \frac{\mu_0 I \hat{z}}{4\pi} \int_{\theta_1}^{\theta_2} \sec \theta d\theta$$

$$\ln |\sec \theta + \tan \theta|$$

(from integral table)

$$\therefore \vec{A} = \frac{\mu_0 I \hat{z}}{4\pi} \ln \left| \underbrace{\frac{\sqrt{(z')^2 + s^2}}{s}}_{\text{positive}} + \underbrace{\frac{z'}{s}}_{\text{positive}} \right| \bigg|_{z=z_1}^{z_2}$$

can drop abs. value
sign

$$\therefore \vec{A} = \frac{\mu_0 I}{4\pi} \left\{ \ln \left(\frac{\sqrt{z_2^2 + s^2}}{s} + \frac{z_2}{s} \right) - \ln \left(\frac{\sqrt{z_1^2 + s^2}}{s} + \frac{z_1}{s} \right) \right\} \hat{z}$$

$$\therefore \vec{A} = \frac{\mu_0 I}{4\pi} \ln \left(\frac{\sqrt{z_2^2 + s^2} + z_2}{\sqrt{z_1^2 + s^2} + z_1} \right) \hat{z}$$

To find $\vec{B}$, we can take $\vec{\nabla} \times \vec{A}$ in cylindrical coords.

$$3. \quad \vec{A}_{in} = \frac{\mu_0 R \omega \sigma}{3} r \sin \theta \hat{\theta}$$

$$\vec{A}_{out} = \frac{\mu_0 R^4 \omega \sigma}{3} \frac{\sin \theta}{r^2} \hat{\theta}$$

$$(a) \quad \vec{B}_{in} = \vec{\nabla} \times \vec{A}_{in} = \frac{1}{r \sin \theta} \left[\frac{\partial}{\partial \theta} (\sin \theta A_{\phi}) \right] \hat{r} - \frac{1}{r} \frac{\partial}{\partial r} (r A_{\theta}) \hat{\theta}$$

$$= \frac{1}{r \sin \theta} \frac{\partial}{\partial \theta} \left(\frac{\mu_0 R \omega \sigma}{3} r \sin^2 \theta \right) \hat{r} - \frac{1}{r} \frac{\partial}{\partial r} \left(\frac{\mu_0 R \omega \sigma}{3} r^2 \sin \theta \right) \hat{\theta}$$

$$= \frac{1}{\cancel{r \sin \theta}} \frac{\mu_0 R \omega \sigma}{3} 2 \cancel{r} \sin \theta \cos \theta \hat{r}$$

$$- \frac{1}{\cancel{r}} \frac{\mu_0 R \omega \sigma}{3} 2 \cancel{r} \sin \theta \hat{\theta}$$

$$= \frac{2 \mu_0 R \omega \sigma}{3} \underbrace{(\cos \theta \hat{r} - \sin \theta \hat{\theta})}_{\hat{z}}$$

$$\vec{B}_{in} = \frac{2}{3} \mu_0 R \omega \sigma \hat{z}$$

In Assign. #4, you will find $\vec{B}$ at the centre of a rotating spherical shell using the Biot-Savart law. You should find this result.

(b) If the solid sphere has charge density $\rho = \frac{Q}{\frac{4}{3}\pi R^3}$,

then the charge on a spherical shell of radius r' & thickness dr' is:

$$\begin{aligned} Q_{shell} &= \rho 4\pi r'^2 dr' \\ &= \frac{Q}{\frac{4}{3}\pi R^3} 4\pi r'^2 dr' \end{aligned}$$

$$\text{s.t. } \sigma = \frac{Q_{shell}}{4\pi r'^2} = \frac{Q dr'}{\frac{4}{3}\pi R^3} = \boxed{\rho dr'}$$

$$\therefore d\vec{A}_{in} = \frac{\mu_0 r' \omega \rho dr'}{3} r \sin \theta \hat{\phi}$$

$$d\vec{A}_{out} = \frac{\mu_0 r'^4 \omega \rho dr'}{3} \frac{\sin \theta}{r^2} \hat{\phi}$$

(c) Add up the shells to construct a sphere of radius R .

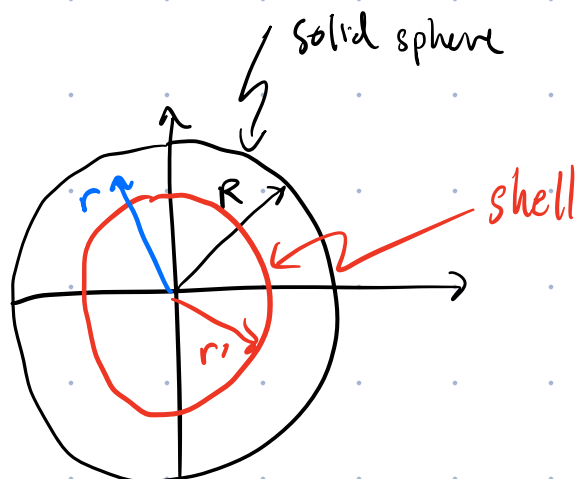
$$\vec{A}_{sphere} = \int_0^r d\vec{A}_{in} + \int_r^R d\vec{A}_{out}$$

$$= \frac{\mu_0 \omega \rho}{3} \sin \theta \hat{\phi} \left[\int_r^R \overbrace{r' r dr'}^{r' > r} + \int_0^r \overbrace{\frac{r'^4}{r^2} dr'}^{r' < r} \right]$$

Cross-sectional view

For $0 < r' < r$,
use $\vec{A}_{out}$

For $r < r' < R$
use $\vec{A}_{in}$



$$\begin{aligned}
&= \frac{\mu_0 \omega \rho}{3} \sin \theta \hat{\phi} \left[\underbrace{\left. \frac{r'^2}{2} r \right|_r^R}_{\frac{R^2 r - r^3}{2}} + \underbrace{\left. \frac{r'^5}{5 r^2} \right|_0^r}_{\frac{r^3}{5}} \right] \\
&\quad \underbrace{\qquad \qquad \qquad}_{\frac{R^2 r}{2} - \frac{(5-2)r^3}{10}} \\
&= r \left(\frac{R^2}{2} - \frac{3}{10} r^2 \right) \\
&= \frac{r}{2} \left(R^2 - \frac{3}{5} r^2 \right)
\end{aligned}$$

$$\therefore \vec{A}_{in} = \frac{\mu_0 \omega \rho}{6} r \left(R^2 - \frac{3}{5} r^2 \right) \sin \theta \hat{\phi}$$

$$(d) \text{ Find } \vec{B}_{in} = \vec{\nabla} \times \vec{A}_{in} = \frac{1}{r \sin \theta} \left[\frac{\partial}{\partial \theta} (\sin \theta A_\phi) \right] \hat{r} - \frac{1}{r} \frac{\partial}{\partial r} (r A_\phi) \hat{\theta}$$

$$\vec{B}_{in} = \frac{\mu_0 \omega \rho}{6} \left\{ \frac{1}{r \sin \theta} \left(\frac{\partial}{\partial \theta} r \left(R^2 - \frac{3}{5} r^2 \right) \sin^2 \theta \right) \hat{r} - \frac{1}{r} \frac{\partial}{\partial r} \left(r^2 \left(R^2 - \frac{3}{5} r^2 \right) \sin \theta \right) \hat{\theta} \right\}$$

$$= \frac{\mu_0 \omega \rho}{6} \left\{ \frac{1}{\cancel{r \sin \theta}} \cancel{r} \left(R^2 - \frac{3}{5} r^2 \right) \cancel{2 \sin \theta} \cos \theta \hat{r} - \frac{1}{\cancel{r}} \left[\cancel{2r} \left(R^2 - \frac{3}{5} r^2 \right) \sin \theta - \frac{6 \cancel{r^2}}{5} \sin \theta \right] \hat{\theta} \right\}$$

$$= \frac{\mu_0 \omega \rho}{6} \left\{ 2 \left(R^2 - \frac{3}{5} r^2 \right) \cos \theta \hat{r} - 2 \left(R^2 - \frac{3}{5} r^2 \right) \sin \theta \hat{\theta} - \frac{6}{5} r^2 \sin \theta \hat{\theta} \right\}$$

$$= \frac{\mu_0 \omega \rho}{6} \left\{ 2 \left(R^2 - \frac{3}{5} r^2 \right) \underbrace{(\cos \theta \hat{r} - \sin \theta \hat{\theta})}_{\hat{z}} - \frac{6}{5} r^2 \sin \theta \hat{\theta} \right\}$$

$$\therefore \vec{B}_{in} = \frac{\mu_0 \omega \rho}{3} \left\{ \left(R^2 - \frac{3}{5} r^2 \right) \hat{z} - \frac{3}{5} r^2 \sin \theta \hat{\theta} \right\}$$

A convenient, but slightly unusual expression since it mixes spherical and Cartesian unit vectors. This time, $\vec{B}_{in}$ is not uniform everywhere inside the sphere.

(e) Finally, at the centre of the rotating sphere $r=0$ s.t.

$$\vec{B}_{in}(0) = \frac{\mu_0 \omega \rho R^2}{3} \hat{z}$$

You will obtain this result in Assign. 4.